The quantum vacuum at the foundation of classical optics

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Maxwell's equations and a

sum rule for elementary particles

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W. Heitler, "The quantum theory of radiation", p.113, 3rd edition, Oxford University 1954

We shall treat the polarization of the vacuum in detail in § 32, but the qualitative consequences are easily seen: A constant field merely leads to a constant polarization. The polarizability is field independent and thus all charges are changed by a universal constant factor (the 'dielectric constant' of the vacuum). Since there is no way of experimenting in an 'ideal vacuum' where this polarizability is absent, this effect is unobservable in principle (although the polarizability turns out to be infinite owing to the infinite number of pairs which contribute). An inhomogeneous field creates, in addition, a non-uniform polarization whose effects are observable (and finite). For example, an additional light wave will be scattered by the polarization dipoles, i.e. in the Coulomb field. Similarly, two light waves will be scattered by each \Box other (§ 32).

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Victor S. Weisskopf: Kongelige Danske Videnskabernes Selskab, Mathematisk-fysiske Meddelelser XIV, No.6, (1936)

Weisskopf's three problematic quantities



These quantities relate to the field free vacuum. It can be taken as self-evident that they can have no physical meaning.

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• Warning

$$c_{
m rel} = \sqrt{rac{E}{m}}$$

 $c_{
m light}$

are not necessarily equivalent! K. Scharnhorst, Phys. Lett. **B 236**, 354(1990)

• We think of the vacuum as a dielectric and diamagnetic medium

 $\mathbf{D} = \varepsilon_0 \mathbf{E} + \mathbf{P} \equiv \mathbf{P}_0 + \mathbf{P}$

$$\mathbf{H} = \frac{1}{\mu_0} \mathbf{B} - \mathbf{M} \equiv \mathbf{M}_0 - \mathbf{M}$$

derive ε_0, μ_0 from properties of the vacuum



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Critical dc electric field strength

$$eE_{
m crit}\lambda_{
m C}=mc_{
m rel}^2$$

Bohr (before 1931), Sauter 1931, Schwinger 1951

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Vacuum impedance determines the power radiated by a dipole



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$$c_{\text{light}} = 1/\sqrt{\varepsilon_0 \mu_0} \qquad c_{\text{light}} = c_{\text{rel}} \sqrt{\frac{2\xi}{\zeta}} \qquad \frac{\zeta}{\xi} = 2$$
$$Z_0 = \frac{2\eta\hbar}{\zeta} \left(\sum_{j}^{\text{e. p.}} q_j^2\right)^{-1} = 8218[\Omega] \frac{\eta}{\zeta} \left(\sum_{j}^{\text{e. p.}} \frac{q_j^2}{e^2}\right)^{-1}$$
$$\sum_{j}^{\text{e. p.}} \frac{q_j^2}{e^2} = \frac{2\hbar}{e^2 Z_0 \zeta} = 21.82 \frac{\eta}{\zeta} \qquad \sum_{j}^{\text{known}} \frac{q_j^2}{e^2} \simeq 10$$

 $\eta\simeq 0.48-2.06$

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$$\alpha = \frac{e^2}{4\pi c\hbar\epsilon_0}$$

$$lpha_0^{-1} = 137.04 = ext{constant} \sum_{j}^{ ext{e. p.}} rac{q_j^2}{e^2} \ lpha_{58\, ext{GeV}}^{-1} = 128.5 \pm 2.5 \ = ext{constant} \sum_{j}^{ ext{e. p.} > 58\, ext{GeV}} rac{q_j^2}{e^2}$$

$$\frac{\alpha_0^{-1}}{\alpha_{58\,\text{GeV}}^{-1}} = \frac{\sum_{\text{all}}}{\sum_{>58\,\text{GeV}}} = \frac{137.04}{128.05 \pm 2.5}$$

$$\sum_{\text{all}} = 104 \left\{ \begin{array}{c} +43 \\ -24 \end{array} \right.$$



Conclusions

- Our simple model predicts a finite number of charged elementary particles and that it relates this number to low-energy properties of the electromagnetic field
- The value predicted by the model is determined by the relative properties of the electric and magnetic interaction of light with the quantum vacuum and is independent of the number of elementary particles.
- We have shown an intimate relationship between the properties of the quantum vacuum and the constants in Maxwell's equations.

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